

Displacement modes on a string

The 1D wave eq. provides general solutions of the form

$$y(t, x) = g\left(t - \frac{x}{v}\right) + h\left(t + \frac{x}{v}\right)$$

where $g(\cdot)$ and $h(\cdot)$ are arbitrary waveforms and v is the wave speed on the string. If we consider one frequency, $\omega = 2\pi f$, we have a general sinusoidal solution:

$$\begin{aligned} y(t, x) &= y^{(+)} \cos\left(\omega\left(t - \frac{x}{v}\right) + \phi^{(+)}\right) + y^{(-)} \cos\left(\omega\left(t + \frac{x}{v}\right) + \phi^{(-)}\right) \\ &= \operatorname{Re}\left\{ y^{(+)} e^{j\left(\omega\left(t - \frac{x}{v}\right) + \phi^{(+)}\right)} + y^{(-)} e^{j\left(\omega\left(t + \frac{x}{v}\right) + \phi^{(-)}\right)} \right\} \end{aligned}$$

where $y^{(+)} = y^{(+)} e^{j\phi^{(+)}}$ and $y^{(-)} = y^{(-)} e^{j\phi^{(-)}}$ are the complex amplitudes of the positive and negative travelling waves, and $y(t, x) = \operatorname{Re}\{y(t, x)\}$ is the real part of the complex displacement field on the string.

At resonance, wave interference creates a standing wave

$$y(t, x) = X(x) e^{j\omega t}$$

To get the mode shape, $X(x)$, we must use the string's boundary conditions, $y_0 = y_L = 0$, which specify how the waves are reflected at the ends.

Mode shapes on a string

In complex form, we have the general sinusoidal behaviour:

$$y(t,x) = y^{(+)} e^{j(\omega(t - \frac{x}{v}))} + y^{(-)} e^{j(\omega(t + \frac{x}{v}))}$$

The first boundary condition at $x=0$ gives

$$y^{(+)} e^{j\omega t} + y^{(-)} e^{j\omega t} = 0$$

$$\Rightarrow y^{(+)} = -y^{(-)}$$

This result simplifies the sinusoidal sound field, via trigonometry:

$$y(t,x) = y^{(+)} (e^{-j\frac{\omega x}{v}} - e^{+j\frac{\omega x}{v}}) e^{j\omega t}$$

$$= -2j y^{(+)} \sin\left(\frac{\omega x}{v}\right) e^{j\omega t}$$

$$\sin x = \frac{e^{jx} - e^{-jx}}{2j}$$

The second boundary condition at $x=L$ gives

$$y^{(+)} e^{j\omega(t - \frac{L}{v})} - y^{(+)} e^{j\omega(t + \frac{L}{v})} = 0$$

$$\Rightarrow y^{(+)} (e^{-j\frac{\omega L}{v}} - e^{+j\frac{\omega L}{v}}) e^{j\omega t} = 0$$

$$\Rightarrow (e^{+j\frac{\omega L}{v}} - e^{-j\frac{\omega L}{v}}) = 0$$

$$\Rightarrow \sin\left(\frac{\omega L}{v}\right) = 0$$

$$\Rightarrow \frac{\omega L}{v} = \{0, \pi, 2\pi, \dots\} = n\pi \text{ for } n = \{0, 1, 2, \dots\}$$

Thus, for resonances with these boundary conditions, we obtain

$$\omega = \frac{n\pi v}{L}$$

and

$$f = \frac{nv}{2L}$$

$$\text{Finally, } y(t,x) = -2j y^{(+)} \overbrace{\sin\left(\frac{n\pi x}{L}\right)}^{x(x)} \cdot e^{j\omega t}$$